



Op-Amps Stability and **Frequency Compensation Techniques**

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Reference: Sergio Franco, "Design with Operational Amplifiers and Analog Integrated Circuits" 4th Edition, Chapter 8, 2015 ■ Harold S. Black, 1927 → Negative feedback concept

Negative feedback provides:

- Gain stabilization
- Reduction of nonlinearity
- Impedance transformation
- But also brings:
 - Potential stability problems
 - Causes accuracy errors for low dc gain
- Here we will discuss frequency compensation techniques

Stability Problem



- Feedback forces x_d to become smaller
- It takes time to detect x_o and feedback to the input
- x_d could be overcorrected (diverge and create instability)
- How to find the optimal (practical) x_d will be based on frequency compensation techniques



Conceptual Gain Margin G_M and Phase Margin ϕ_m G_M is the number of dBs by which $|T(j\omega_{-180^\circ})|$ can increase until it becomes 0 dB

$$G_M = 20\log\frac{1}{|T(j\omega_{-180^\circ})|}$$

Phase margin ϕ_m is the number of degrees by which $\angle T(j\omega_x)$ can be reduced until it reaches $-\pi$ (-180°) $\phi_m = 180^\circ + \angle T(j\omega_x)$

or

$$\angle T(j\omega_x) = \phi_m - 180^\circ$$

* ω_x is the crossover frequency

At the crossover point,

$$T(j\omega_x) = 1 \cdot \angle T(j\omega_x) = 1 \cdot \angle (\phi_m - 180^\circ) = -e^{j\phi_m}$$

The non-ideal closed loop transfer function becomes

$$H_{CL}(j\omega_{x}) = \frac{A(j\omega_{x})}{1+\beta A(j\omega_{x})} = \frac{A(j\omega_{x})}{1+T(j\omega_{x})} = \frac{A_{ideal}}{1+1/T(j\omega_{x})}$$
$$= \frac{A_{ideal}}{1-e^{-j\omega_{x}}} = \frac{A_{ideal}}{1-(\cos\phi_{m}-j\sin\phi_{m})}$$
$$|H_{CL}(j\omega_{x})| = |A_{Ideal}| \frac{1}{\sqrt{(1-\cos\phi_{m})^{2}+\sin^{2}\phi_{m}}}, A_{Ideal} = \frac{1}{\beta}$$

• Observe that different ϕ_m yield different errors. i.e.

ϕ_m	$ H_{CL}(j\omega_x) $
90°	0.707
60°	1.00
45°	1.31
30°	1.93
15°	3.83
0°	∞ (oscillatory behavior)

- In practical systems, $\phi_m = 60^\circ$ is required
- A worst case $\phi_m = 45^{\circ}$ for a typical lower limit
- For $\phi_m < 60^\circ$, we have $|A(j\omega_x)| > |A_{ideal}|$

indicating a peaked closed-loop response.

Why a dominant pole is required for a stable amplifier?

A good Op amp design implies:

(i) $|\beta A(j\omega)| >> 1$ over as wide a band of frequencies as possible (ii) The zeroes of $\beta A(j\omega) - 1 = 0$ must be all in the left-hand plane Note that $\beta A(j\omega) < 0$

These two conditions often conflict with each other. These trade-offs should be carefully considered. Let's consider a practical amplifier characterized with these poles,

$$A(s) = -\frac{A_0}{(1 + s / \alpha_1)(1 + s / \alpha_2)(1 + s / \alpha_3)} = \frac{-A_0 \alpha_1 \alpha_2 \alpha_3}{(s + \alpha_1)(s + \alpha_2)(s + \alpha_3)}$$

The characteristic equation becomes

$$\beta A(s) - 1 = \frac{-\beta A_0 \alpha_1 \alpha_2 \alpha_3}{(s + \alpha_1)(s + \alpha_2)(s + \alpha_3)} - 1 = 0$$
$$(s + \alpha_1)(s + \alpha_2)(s + \alpha_3) + \beta A_0 \alpha_1 \alpha_2 \alpha_3 = 0$$

 $s^{3} + s^{2}(\alpha_{1} + \alpha_{2} + \alpha_{1}) + s(\alpha_{1}\alpha_{2} + \alpha_{1}\alpha_{3} + \alpha_{2}\alpha_{3}) + \alpha_{1}\alpha_{2}\alpha_{3}(1 + \beta A_{0}) = 0$

Note that βA_0 is the critical parameter that determines the pole locations for a given α_1 , α_2 and α_3 ($0 \le \beta \le 1$). Furthermore, when $\beta A_0 = 0$, the roots are at $-\alpha_1$, $-\alpha_2$ and $-\alpha_3$. Therefore, for small βA_0 , the roots should be in the left-hand plane (LHP). However, for $\beta A_0 >> 1$, two of the roots might be forced to move to the righthand plane (RHP). This can be verified by applying Routh's stability criterion. Let us write the polynomial as

$$b_3s^3 + b_2s^2 + b_1s + b_0 = 0$$

In order to have, in the above equation, left half plane roots, all the coefficients must be positive and satisfy

$$b_2 b_1 - b_3 b_0 = 0$$

The condition for imaginary-axis roots become

$$b_{3}(j\omega)^{3} + b_{2}(j\omega)^{2} + b_{1}(j\omega) + b_{0} = 0$$

$$(b_{0} - b_{2}\omega^{2}) + j\omega(b_{1} - b_{3}\omega^{2}) = 0$$

Now, for $s=j\omega$ being a root, both real and imaginary parts must be zero. That is,

$$b_0 - b_2 \omega^2 = 0$$
, $b_1 - b_3 \omega^2 = 0$ or $b_3 b_0 = b_1 b_2$

Then the two roots are placed at

$$\omega_{p_{2,3}} = \pm j \sqrt{\frac{b_0}{b_2}} = \pm j \sqrt{\frac{b_1}{b_3}}$$

Critical Value of βA_{θ}

$$b_0 = \frac{b_2 b_1}{b_3}$$

Then

$$\alpha_1 \alpha_2 \alpha_3 \left[1 + \left(\beta A_0 \right)_C \right] = \left(\alpha_1 + \alpha_2 + \alpha_3 \right) \left(\alpha_1 \alpha_2 + \alpha_1 \alpha_3 + \alpha_2 \alpha_3 \right)$$

Thus, the critical value of βA_0 becomes

$$\left(\beta A_0\right)_C = 2 + \frac{\alpha_1}{\alpha_2} + \frac{\alpha_1}{\alpha_3} + \frac{\alpha_2}{\alpha_1} + \frac{\alpha_2}{\alpha_3} + \frac{\alpha_3}{\alpha_1} + \frac{\alpha_3}{\alpha_2}\right)_C$$

Thus when βA_0 becomes $(\beta A_0)_C$, the amplifier will oscillate at

$$\omega_{OSC} = \sqrt{\frac{b_1}{b_3}} = \sqrt{\alpha_1 \alpha_2 + \alpha_1 \alpha_3 + \alpha_2 \alpha_3}$$

Also when $\beta A_0 > (\beta A_0)_C$, the amplifier has RHP poles, therefore is **unstable**.

Critical Value of βA_{θ}

Let us consider some numerical examples. Let $A_0 = 10^5$, (i) Three equal poles $\alpha_1 = \alpha_2 = \alpha_3 = 10^7 \text{ rad/s}$. The amplifier oscillates at $\omega_{OSC} = \alpha_1 \sqrt{3} = 10^7 \sqrt{3} \text{ rad/s}$

$$(\beta A_0)_C = 8, \quad \beta_C = 8 / A_0 = 8 \times 10^{-5}$$

(ii)
$$\alpha_1 = \frac{\alpha_2}{10^4} = \frac{\alpha_3}{10^4}$$
, then the critical loop gains yields
 $(\beta A_0)_C \cong 2\frac{\alpha_2}{\alpha_1} = 2 \times 10^4$, thus the amplifier is stable if $\beta_C < \frac{2 \times 10^4}{A_0} = 0.2$

Since $\beta = R_1 / (R_1 + R_2)$, $\beta < 0.2$ causes $(R_2 / R_1) > 4$, which means that for an inverting (non-inverting) configuration,

the gain must be greater than -4(5) to keep the amplifier stable.

(iii) Let us determine A_{0C} under the most stringent condition $\beta=1$. Then from previous equation,

$$A_0 < 2 + \frac{\alpha_1}{\alpha_2} + \frac{\alpha_1}{\alpha_3} + \frac{\alpha_2}{\alpha_1} + \frac{\alpha_2}{\alpha_3} + \frac{\alpha_3}{\alpha_1} + \frac{\alpha_3}{\alpha_2}$$

In order to have a large A_0 , the poles must be widely separated. i.e. $\alpha_1 << \alpha_2 << \alpha_3$, then the A_0 inequality can be approximated as

$$A_0 < \frac{\alpha_2}{\alpha_1} + \frac{\alpha_3}{\alpha_1}$$

To obtain a conservative A_0 , let $\alpha_2 = \alpha_3$, which yields

$$A_0 < 2\frac{\alpha_2}{\alpha_1}$$

This inequality bounds the DC gain to provide a stable closed loop configuration.

- Peaking in the frequency domain usually implies ringing in the time domain
- Normalized second-order all-pole (low pass) system

$$H(s)\Big|_{s=j\omega} = \frac{\omega_0^2}{s^2 + \frac{\omega_0}{Q}s + \omega_0^2}\Big|_{s=j\omega} = \frac{1}{1 - \frac{\omega^2}{\omega_0^2} + j\frac{1}{Q}\cdot\frac{\omega}{\omega_0}}$$



■ *GP*: Peak gain – We have the error function,

$$\left|E(j\omega)\right| = \frac{1}{\sqrt{\left(1 - \frac{\omega^2}{\omega_0^2}\right)^2 + \frac{1}{Q^2} \cdot \frac{\omega^2}{\omega_0^2}}} = \frac{1}{\sqrt{D(\omega)}} \tag{1}$$

To find out the maximum value of $|E(j\omega)|$, calculate the derivative of $D(\omega)$ and make it equal to zero

$$\frac{d}{d\omega}D(\omega) = 2\left(\frac{1}{Q^2} - 2\right)\frac{\omega}{\omega_0^2} + \frac{4\omega^3}{\omega_0^4} = 0 \quad \implies \quad \omega_*^2 = \left(1 - \frac{1}{2Q^2}\right)\omega_0^2 \quad \text{OR} \quad \omega = 0$$

For $Q > 1/\sqrt{2}$, use ω_* in Eq. (1) and we get the peak gain

$$GP = \frac{2Q^2}{\sqrt{4Q^2 - 1}} \ge \left| E(j0) \right| = 1$$
⁽²⁾

• **OS:** Overshoot – Inverse Laplace transform $s - \text{domain} : \frac{b}{(s+a)^2 + b^2} \xrightarrow{\text{Inverse}\\\text{Laplace}} \text{time} - \text{doamin} : e^{-at} \sin(bt)u(t)$

For a 2nd order all pole error function, the **impulse response** is

$$H(s)\Big|_{s=j\omega} = \frac{\omega_0^2}{s^2 + \frac{\omega_0}{Q}s + \omega_0^2} = \frac{\omega_0^2}{b} \cdot \frac{b}{(s+a)^2 + b^2} \qquad v_{in}(t) = \delta(t)$$

$$a = \frac{\omega_0}{2Q}$$

$$\xrightarrow{Inverse}_{Laplace} \to h(t) = \frac{\omega_0^2}{b}e^{-at}\sin(bt)u(t) \qquad b = \sqrt{1 - \frac{1}{4Q^2}\omega_0}$$

Consider the normalized step response of this system,

$$y(t) = \int_0^t h(t)dt = 1 - \frac{\omega_0}{b} e^{-at} \sin(bt + \phi), \phi = \tan^{-1}\frac{b}{a} \quad \text{for } v_{in}(t) = u(t) = \int \delta(t)dt \quad (3)$$

Use the damping factor ξ to represent, $a = \xi \omega_0, b = \sqrt{1 - \xi^2} \omega_0$

Usually, the damping factor $\xi=1/2Q$ is used to characterize a physical 2nd order system. Thus, we rewrite the normalized time-domain equation (3)

$$y(t) = 1 - \frac{\omega_0}{b} e^{-at} \sin(bt + \phi) = 1 - \frac{1}{\sqrt{1 - \xi^2}} e^{-\xi\omega_0 t} \sin\left(\sqrt{1 - \xi^2}\omega_0 t + \phi\right)$$

$$\phi = \tan^{-1}\frac{b}{a} = \tan^{-1}\frac{\sqrt{1 - \xi^2}}{\xi}$$

$$\phi \approx \frac{\sqrt{1 - \xi^2}}{\xi}, \text{ for small } \xi$$

$$(4)$$

For under damped case $\xi < 1$ (Q > 0.5), the overshoot is the peak value of y(t). To find the peak value, we first calculate the derivative of y(t), and make it equal to 0.

$$\frac{d}{d(\omega_0 t)} y(t) = -\frac{\xi}{\sqrt{1-\xi^2}} e^{-\xi\omega_0 t} \sin\left(\sqrt{1-\xi^2}\omega_0 t + \phi\right) + e^{-\xi\omega_0 t} \cos\left(\sqrt{1-\xi^2}\omega_0 t + \phi\right)$$

$$\frac{d}{d(\omega_0 t)} y(t) = 0$$

$$\Rightarrow \frac{\xi}{\sqrt{1 - \xi^2}} e^{-\xi\omega_0 t} \sin\left(\sqrt{1 - \xi^2}\omega_0 t + \phi\right) = e^{-\xi\omega_0 t} \cos\left(\sqrt{1 - \xi^2}\omega_0 t + \phi\right)$$

$$\Rightarrow \tan\left(\sqrt{1 - \xi^2}\omega_0 t + \phi\right) = \frac{\sqrt{1 - \xi^2}}{\xi} = \tan\phi$$

To satisfy the equality above, we have

$$\omega_0 t = \frac{n\pi}{\sqrt{1-\xi^2}}, n = 0, 1, 2, \dots$$

In other words, the **normalized step response** y(t) in Eq.(3) achieves **extreme values** at time steps of n=0, 1, 2, ...



$$n = 0, \quad \omega_0 t = 0, \qquad y(t) = 0$$

$$n = 1, \quad \omega_0 t = \frac{\pi}{\sqrt{1 - \xi^2}}, \quad y(t) = 1 - \frac{1}{\sqrt{1 - \xi^2}} e^{-\frac{\pi\xi}{\sqrt{1 - \xi^2}}} \sin(\pi + \phi) = 1 + e^{-\frac{\pi\xi}{\sqrt{1 - \xi^2}}}$$

$$n = 2, \quad \omega_0 t = \frac{2\pi}{\sqrt{1 - \xi^2}}, \quad y(t) = 1 - \frac{1}{\sqrt{1 - \xi^2}} e^{-\frac{2\pi\xi}{\sqrt{1 - \xi^2}}} \sin(2\pi + \phi) = 1 - e^{-\frac{2\pi\xi}{\sqrt{1 - \xi^2}}}$$

$$n = 3, \quad \omega_0 t = \frac{3\pi}{\sqrt{1 - \xi^2}}, \quad y(t) = 1 - \frac{1}{\sqrt{1 - \xi^2}} e^{-\frac{3\pi\xi}{\sqrt{1 - \xi^2}}} \sin(3\pi + \phi) = 1 + e^{-\frac{3\pi\xi}{\sqrt{1 - \xi^2}}}$$

$$\dots \dots \dots \dots \dots \dots$$

Therefore, the global peak value of y(t) is achieved when n=1.

Thus, the overshoot is defined as

$$OS(\%) = 100 \times \frac{\text{Peak Value} - \text{Final Value}}{\text{Final Value}} = 100 \times e^{-\frac{\pi\xi}{\sqrt{1-\xi^2}}}$$
(5)

• ϕ_m : Phase Margin For a 2nd order all-pole error function $F(s) = \frac{T(s)}{1} = 1$

$$E(s) = \frac{T(s)}{1 + T(s)} = \frac{1}{1 + 1/T(s)} = \frac{1}{\frac{s^2}{\omega_0^2} + \frac{1}{Q}\frac{s}{\omega_0} + 1}$$

Therefore, the loop gain T(s) is given by

$$T(s) = \frac{1}{\frac{s^2}{\omega_0^2} + \frac{1}{Q}\frac{s}{\omega_0}}$$

The cross-over frequency thus can be obtained

$$\left|T(j\omega_{x})\right| = \left(\sqrt{\left(\frac{\omega_{x}^{2}}{\omega_{0}^{2}}\right)^{2} + \left(\frac{\omega_{x}}{\omega_{0}Q}\right)^{2}}\right)^{-1} = 1$$

Solve the equation and get the crossover frequency,

$$\omega_{x} = \omega_{0} \left(\sqrt{\frac{1}{4Q^{4}} + 1} - \frac{1}{2Q^{2}} \right)^{1/2} = \omega_{0} \left(\sqrt{4\xi^{4} + 1} - 2\xi^{2} \right)^{1/2}$$

And thus the phase margin is

$$\phi_m = 180^\circ + \angle T(j\omega_x) = \cos^{-1}\left(\sqrt{1 + \frac{1}{4Q^4}} - \frac{1}{2Q^2}\right) = \cos^{-1}\left(\sqrt{4\xi^4 + 1} - 2\xi^2\right)$$

Study the relationship between phase margin and gain peaking (Eq.2) or overshoot (Eq.5), we have

$$GP(60^{\circ}) \cong 0.3 dB$$
 $OS(60^{\circ}) \cong 8.8\%$ $Q \approx 0.82$
 $GP(45^{\circ}) \cong 2.4 dB$ $OS(45^{\circ}) \cong 23\%$ $Q \approx 1.18$



One effective method of assessing stability for minimum phase systems



from the magnitude Bode plots is by determining the ROC.

The Rate of Closure (ROC)

Determining the ROC is done by observing the slopes of |A|and $|1/\beta|$ at their intersection point (cross-over frequency f_x) and deciding the magnitude of their difference.

$$ROC = |Slope(|A|) - Slope(|1/\beta|)_{f=f_x}$$

The ROC is used to estimate the phase margin and therefore the stability (How ?)

The Rate of Closure (ROC)

Observing a single-root transfer function



This correlation holds also if H(s) has more than one root, provided the roots are real negative, and well separated, say, at least a decade apart.

The Rate of Closure (ROC)

In a feedback system, suppose both |A|and $|1/\beta|$ have been graphed.

$$\angle T(jf_x) = \angle A(jf_x) - \angle \beta^{-1}(jf_x)$$
$$\cong -4.5 \times ROC$$



Thus, the ROC can be used to estimate the phase margin (Page 4)

$$\phi_m = 180^\circ + \angle T(jf_x)$$

Cases

$$ROC \cong 20 dB / dec \implies \phi_m \cong 90^\circ$$
$$ROC \cong 30 dB / dec \implies \phi_m \cong 45^\circ$$
$$ROC \cong 40 dB / dec \implies \phi_m \cong 0^\circ$$
$$ROC > 40 dB / dec \implies \phi_m < 0^\circ$$

The Rate of Closure (ROC)



Stability in Constant-GBP OpAmp

Constant-GBP OpAmp (i.e. $A(s)|_{s=j\omega} = \frac{\omega_t}{j\omega}$)

- Unconditionally stable with frequency-independent feedback, or $\angle \beta = 0$. (e.g. in a non-inverting or inverting amplifier, the feedback network contains only resistors)
- Stable for any $\beta \leq 1$.
- In feedback systems, since now we have $(T, f(x)) = (f(x))^{-1}$

$$\angle T = \angle (A\beta) = \angle A$$
, $\angle A \cong -90^\circ$

these circuits enjoy

$$\phi_m = 180^\circ + \angle A(jf_x) \cong 180^\circ - 90^\circ \cong 90^\circ$$

• Typically, due to additional high-order poles in OpAmps,

$$60^\circ \!\leq\! \phi_m \!\leq\! 90^\circ$$

Feedback Pole

Feedback Pole

Feedback network includes reactive elements \rightarrow Stability may no longer be unconditional



The effect of a pole within the feedback loop

Feedback Pole

Examine the error function

$$E(s) = \frac{H_{CL}(s)}{A_{ideal}} = \frac{1}{1+1/T}, \qquad T = A\beta, \quad A_{ideal} = \frac{1}{\beta}$$

Using OpAmp high-frequency approximation: $A(j\omega) = \frac{GB}{j\omega} = \frac{f_t}{if}$

$$E(s) = \frac{1}{1 + \frac{1}{A\beta}} = \frac{1}{1 + j\frac{f}{f_t\beta_0} - \frac{f^2}{f_t\beta_0 f_z}}$$

Refer to page 14, and we have $s=j2\pi f$. The peak value of E(s) can be obtained. For Q >> 1, the approximate result is

$$f_x = \sqrt{f_z \beta_0 f_t}, Q = \sqrt{\beta_0 f_t / f_z}$$

Feedback Pole

The lower f_z compared to $\beta_0 f_t$, the higher the Q and, hence, the more pronounced the peaking and ringing.

Derive the phase margin

$$\angle T(jf_x) = \angle A(jf_x) - \angle |1/\beta(jf_x)| \cong -90^\circ - \tan^{-1}(f_x/f_z)$$
$$\frac{f_x}{f_z} = \sqrt{\frac{\beta_0 f_t}{f_z}}$$

As $f_z \ll \beta_0 f_t$, $\angle T(jf_x) \cong -180^\circ$ and ROC = 40dB/dec The circuit is on the verge of oscillation!

Differentiator

- Feedback pole example: differentiator
- Assume constant-GBP OpAmp $A(jf) \cong GB / j\omega \cong f_t / jf$ $\beta = \frac{Z_C}{Z_C + R} = \frac{1}{1 + jf / f_z}, Z_C = \frac{1}{j\omega C}$
- To stabilize the differentiator, add a series resistance Rs.



Differentiator

- At low frequency, R_S has little effect because $R_S << |1/j\omega C|$
- At high frequency, C acts as a short compared to $R_{S_{,}}$ The feedback network becomes $|1/\beta| = 1 + R/R_{S}$.



Differentiator

Assume $R_S \ll R_C$, the series resistor R_S introduces an extra pole frequency f_e

$$\frac{1}{\beta} = \frac{Z_C + R}{Z_C} \approx \frac{1}{\beta_0} \frac{1 + jf / f_z}{1 + jf / f_e} \qquad \qquad Z_C = \frac{1}{j\omega C} + R_S, \beta_0 = 1, f_x = \sqrt{f_t f_z}$$

• Choose $R_s \approx R / \sqrt{f_t / f_z}$, we have $f_e = \sqrt{f_t f_z}$

$$\angle T(jf_x) = \angle A(jf_x) - \angle |1/\beta(jf_x)|$$

$$\cong -90^\circ + \tan^{-1}(f_x/f_e) - \tan^{-1}(f_x/f_z)$$

$$= -135^\circ$$

Therefore, ROC = 30 dB/dec, $\phi_m \approx 45^{\circ}$

Stray Input Capacitance Compensation

All practical OpAmps exhibit stray input capacitance. The net capacitance C_n of the inverting input toward ground is

$$C_n = C_d + C_C / 2 + C_{ext}$$

 C_d is the differential Cap between input pins, $C_c/2$ is the common-mode cap of each input to ground, and C_{ext} is the external parasitic cap.

In the absence of C_f , there's a pole in feedback $\frac{1}{\beta} = (1 + R_2 / R_1) \{1 + jf [2\pi (R_1 / / R_2)C_n]\}$ ROC ≈ 40 dB/dec (See page 29)



Stray Input Capacitance Compensation

- Solution: Introduce a feedback capacitance C_f to create feedback phase lead.
- In the presence of C_f we have

$$\frac{1}{\beta} = \left(1 + \frac{R_2}{R_1}\right) \frac{1 + jf/f_z}{1 + jf/f_p}$$
$$f_z = \frac{1}{2\pi (C_n + C_f)(R_1 \| R_2)} \quad , \quad f_p = \frac{1}{2\pi C_f R_2}$$

• To have $\phi_m = 45^\circ$ (i.e. ROC=30 dB/dec):

Make the cross-over frequency exactly at f_p

$$|A(jf_p)| = \left|\frac{1}{\beta(jf_p)}\right| \cong \frac{1}{\beta_{\infty}} \to |A(jf_p)| \cong 1 + \frac{C_n}{C_f}$$

Since
$$|A(jf_p)| = \frac{f_t}{f_p} \to \frac{1}{f_p} = 2\pi C_f R_2 = \frac{1}{f_t} \left(1 + \frac{C_n}{C_f}\right)$$

Solve the equation to get C_f
$$1 + R_2/R_1$$

$$C_f = \frac{1 + \sqrt{1 + 8\pi R_2 C_n f_t}}{4\pi R_2 f_t}$$



Stray Input Capacitance Compensation

• To have $\phi_m = 90^\circ$ (i.e. ROC=20 dB/dec):

Place f_p exactly on the top of f_z to cause a pole-zero cancellation

$$f_z = f_p$$
$$(C_n + C_f)(R_1 || R_2) = C_f R_2$$

Thus using simple algebra




Capacitive-Load Isolation

- There're applications in which the external load is heavily capacitive.
- Load capacitance C_L
 - A new pole is formed with output resistance r_o and C_L
 - Ignore loading by the feedback network



- The loaded gain is $A_{loaded} \cong A(1 + jf / f_p)^{-1}, f_p = (2\pi r_o C_L)^{-1}$
- ROC is increased and thus invite instability
- Solution: Add a small series resistance R_S to decouple the output from C_L

Capacitive-Load Isolation



$$R_{S} = \frac{R_{1}}{R_{2}}r_{O} \qquad C_{f} = \left(1 + \frac{R_{1}}{R_{2}}\right)^{2} \left(\frac{r_{o}}{R_{2}}\right)C_{L}$$

- The poles of the uncompensated OpAmp are located close together thus accumulating about 180° of phase shift before the 0-dB crossover frequency f_x .
- Unstable device, thus efforts must be done to stabilize it.
- Example of uncompensated OpAmps is 748, which is the uncompensated version of 741.

They can be approximated as a three-pole system

$$a(jf) = \frac{a_0}{(1 + jf/f_1)(1 + jf/f_2)(1 + jf/f_3)}$$



Three-pole OpAmp model

Stability of Uncompensated OpAmp

• With a frequency-independent feedback (i.e. $1/\beta$ curve is flat) around an uncompensated OpAmp, we have

 $|T| = |a|\beta$

|T| curve can be visualized as the |a| curve with the $1/\beta$ line is the new 0-dB axis

• For $1/\beta \ge |a(jf_{-135^\circ})|$ ROC $\le 30 \text{ dB/dec}$

 $\phi_m \ge 45^\circ$

 $\phi_m < 0^\circ$

- For $|a(jf_{-180^{\circ}})| \le 1/\beta < |a(jf_{-135^{\circ}})|^{|a(jf_{-135^{\circ}})|}$ $30 \text{ dB/dec} \le \text{ROC} \le 40 \text{ dB/dec}$ $0^{\circ} \le \phi_m \le 45^{\circ}$
- For $1/\beta \le |a(jf_{-180^\circ})|$ ROC < 40 dB/dec



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Stability of Uncompensated OpAmp

- The uncompensated OpAmp provides only adequate phase margin only in high-gain applications (i.e. high $1/\beta$)
- To provide adequate phase margin in low-gain application, frequency compensation is needed.
 - Internal compensation \rightarrow Achieved by changing a(jf)
 - External compensation \rightarrow Achieved by changing $\beta(jf)$

Internal Frequency Compensation

- How to stabilize the circuit by modifying the open loop response a(jf)?
 - Dominant-Pole Compensation
 - Shunt-Capacitance Compensation
 - Miller Compensation
 - Pole-Zero Compensation
 - Feedforward Compensation

Dominant-Pole Compensation

An additional pole at sufficiently low frequency is created to insure a roll-off rate of -20 dB/dec all the way up to the crossover frequency.
Additional Pole



Dominant-Pole Compensation

Numerical example:

- $r_d = \infty, r_o = 0$
- $g_1=2 \text{ mA/V}, R_1=100 \text{ k}\Omega, g_2=10 \text{ mA/V}, R_2=50 \text{ k}\Omega$
- $f_1 = 100 \ kHz, f_2 = 1 \ MHz, f_3 = 10 \ MHz$
- Find the required value of f_d for $\phi_m = 45^\circ$ with $\beta = 1$

For $\phi_m = 45^\circ$, we have

$$f_x = f_1$$

Draw a straight line of slope -20 dB/dec until it intercepts with the DC gain asymptote at point D and get f_d .

$$\frac{|a(jf_{x(new)})|}{|a(jf_d)|} = \frac{f_d}{f_{x(new)}}$$

Thus:

$$f_d = \frac{f_{x(new)}}{\beta a_0} = \frac{f_1}{\beta \times (g_1 R_1 g_2 R_2)} = 1 Hz$$

Requires EXTREMELY LARGE passive components



Shunt-Capacitance Compensation

- The dominant-pole technique adds a fourth pole \rightarrow Extra cost and less bandwidth.
- This technique rearranges the existing rather than creating a new pole.
- It decreases the first (dominant) pole to sufficiently low frequency to insure a roll-off rate of -20 dB/dec all the way up to the crossover frequency.



Shunt-Capacitance Compensation



Since f_{1(new)} is chosen to insure roll-off rate of -20 dB/dec all the way up to the crossover frequency.
Thuse

Thus:

The first pole is decreased by adding an extra capacitance to the internal node causing it.



Given the value of $f_{1(new)}$ from the desired value of ϕ_m , we can find C_c

$$f_{1(new)} = \frac{f_x}{\beta a_0} = \frac{1}{2\pi R_1 (C_1 + C_c)} \quad \rightarrow \quad C_c \cong \frac{\beta a_0}{2\pi R_1 f_x}$$

Shunt-Capacitance Compensation

Numerical example:

- $r_d = \infty, r_o = 0$
- $g_1 = 2 \text{ mA/V}, R_1 = 100 \text{ k}\Omega$
- $g_2 = 10 \text{ mA/V}, R_2 = 50 \text{ k}\Omega$
- $f_1 = 100 \ kHz$, $f_2 = 1 \ MHz$, $f_3 = 10 \ MHz$
- Find the required value of C_c for $\phi_m = 45^\circ$ with $\beta = 1$



For
$$\phi_m = 45^\circ$$
, $f_x = f_2 = 1 MHz$
Then, $f_{1(new)} = \frac{f_2}{a_0\beta} = \frac{f_2}{g_1R_1g_2R_2} = 10 Hz \rightarrow C_c = \frac{1}{2\pi R_1f_{1(new)}} = 159 nf$

EXTREMELY LARGE Unsuitable for monolithic fabrication

Miller's Theorem

If A_v is the voltage gain from node 1 to 2, then a floating impedance Z_F can be converted to two grounded impedances Z_1 and Z_2 :



Miller Compensation

Applying Miller's theorem to a floating capacitance connected between the input and output nodes of an amplifier.



- The floating capacitance is converted to two grounded capacitances at the input and output of the amplifier.
- The capacitance at the input node is larger than the original floating capacitance (Miller multiplication effect)

This technique places a capacitor C_c in the feedback path of one of the internal stages to take advantage of Miller multiplication of capacitors.



The reflected capacitances due to C_c and the DC voltage gain between V_2 and V_1 ($a_2 = -g_2R_2$) yields

$$C_{1,c} = C_c (1 + g_2 R_2) \quad \text{and} \quad C_{2,c} = C_c \left(1 + \frac{1}{g_2 R_2} \right)$$
$$\cong |a_2|C_c \quad \cong C_c$$

A low-frequency dominant pole can be created with a moderate capacitor value.

Miller Compensation

Accurate transfer function

$$\frac{V_2}{V_d} \cong g_1 R_1 g_2 R_2 \frac{1 - jf/f_z}{(1 - jf/f_{1(new)})(1 - jf/f_{2(new)})}$$



Pole/zero locations

$$\begin{split} \omega_{z} &= \frac{g_{2}}{c_{c}} & \text{RHP zero} \\ \omega_{1(new)} &= \frac{1}{R_{1}C_{1} + g_{2}R_{2}R_{1}C_{c} + R_{2}C_{2}} \cong \frac{1}{R_{1}g_{2}R_{2}C_{c}} = \frac{1}{|a_{2}|C_{c}R_{1}} & \text{Dominant Pole} \\ \omega_{2(new)} &= \frac{R_{1}C_{1} + R_{1}g_{2}R_{2}C_{c} + R_{2}C_{2}}{R_{1}R_{2}(C_{1}C_{c} + C_{1}C_{2} + C_{c}C_{2})} \cong \frac{g_{2}C_{c}}{C_{1}C_{c} + C_{1}C_{2} + C_{c}C_{2}} & \text{Second Pole} \end{split}$$

Right-half plane zero:

• The RHP zero is a result of the feedforward path through C_c



- The circuit is no longer a minimum-phase system.
- It introduces excessive phase shift, thus reduces the phase margin.
- In bipolar OpAmps, it is usually at much higher frequency than the poles $\rightarrow 1 f/f_z \cong 1$

Miller Compensation

Pole Splitting

- Increasing C_c lowers $f_{1(new)}$ and raises $f_{2(new)}$
- The shift in f_2 eases the amount of shift required by $f_1 \rightarrow$ Higher bandwidth



• Increasing C_c above a certain limit makes f_2 stops to increase.

$$\omega_{2(new)} = 2\pi f_{2(new)} = \frac{g_2}{C_1 + \frac{C_1 C_2}{C_c} + C_2} \cong \frac{g_2}{C_1 + C_2} \bigg|_{C_c \gg C_1, C_2}$$

Miller Compensation

Numerical example:

- $r_d = \infty$, $r_o = 0$, $g_1 = 2 \text{ mA/V}$, $R_1 = 100 \text{ k}\Omega$, $g_2 = 10 \text{ mA/V}$, $R_2 = 50 \text{ k}\Omega$
- $f_1 = 100 \ kHz, f_2 = 1 \ MHz, f_3 = 10 \ MHz$
- Find the required value of C_c for $\phi_m = 45^\circ$ with $\beta = 1$



From f_1 and f_2 , we can calculate $C_1 = 15.9 \ pF$ and $C_2 = 3.18 \ pF$ Assume C_c is large $\rightarrow f_{2(new)} = \frac{g_2}{2\pi(C_1+C_2)} = 83.3 \ MHz > f_3$ Since $f_{2(new)} > f_3 \rightarrow f_3$ is the first non-dominant pole For $\phi_m = 45^\circ$, $f_x = f_3 = 10 \ MHz$ Then, $f_{1(new)} = \frac{f_3}{a_0\beta} = \frac{f_3}{g_1R_1g_2R_2} = 100 \ Hz$ $C_c = \frac{1}{2\pi R_1g_2R_2f_{1(new)}} = 31.8 \ pF$ Much smaller than that of shuntcapacitance compensation Suitable for monolithic fabrication This technique uses a large compensation capacitor ($C_c \gg C_1$) to lower the first pole f_1 .

It also uses a small resistor $(R_c \ll R_1)$



to create a zero that cancels the second pole f_2 .

- The compensated response is then dominated by the lowered first pole up to f_3
- Transfer function

$$\frac{V_1}{V_d} = -g_1 R_1 \frac{1 + jf/f_z}{\left(1 + jf/f_{1(new)}\right)\left(1 + jf/f_4\right)}$$

Pole/zero locations

$$f_{1(new)} \cong \frac{1}{2\pi R_1 C_c}$$
 , $f_z = \frac{1}{2\pi R_c C_c}$, $f_4 \cong \frac{1}{2\pi R_c C_1}$

- C_c and R_c lowers the dominant pole $f_{1(new)} \ll f_1$, creates a zero f_z , and creates an additional pole $f_4 \gg f_z$
- Choose R_c such that f_z cancels f_2
- The open loop gain now becomes

$$a_{new}(jf) = \frac{a_0}{\left(1 + jf/f_{1(new)}\right)\left(1 + jf/f_3\right)\left(1 + jf/f_4\right)}$$



- To have $\phi_m = 45^\circ$, the cross-over frequency should be f_3
 - Since the compensated response is dominated by $f_{1(new)}$ pole up to f_3 $\frac{|a(jf_{1(new)})|}{|a(jf_3)|} = \frac{a_0}{1/\beta} = \frac{f_3}{f_{1(new)}}$

Thus,
$$f_{1(new)} = f_3/a_0\beta$$

Numerical example:

- $r_d = \infty$, $r_o = 0$, $g_1 = 2 \text{ mA/V}$, $R_1 = 100 \text{ k}\Omega$, $g_2 = 10 \text{ mA/V}$, $R_2 = 50 \text{ k}\Omega$
- $f_1 = 100 \ kHz, f_2 = 1 \ MHz, f_3 = 10 \ MHz$
- Find the required value of C_c for $\phi_m = 45^\circ$ with $\beta = 1$

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From f_1 and f_2 , we can calculate $C_1 = 15.9 \ pF$ and $C_2 = 3.18 \ pF$ For $\phi_m = 45^\circ$, $f_x = f_3 = 10 \ MHz$ Then, $f_{1(new)} = \frac{f_3}{a_0\beta} = \frac{f_3}{g_1R_1g_2R_2} = 100 \ Hz \rightarrow C_c = \frac{1}{2\pi R_1f_{1(new)}} = \underbrace{15.9 \ nf}_{Relaxed \ compared \ to \ shunt-capacitance \ compared \ to \ shunt-capacitance \ compensation, \ but \ still \ LARGE$ Since $f_4 = \frac{1}{2\pi R_c C_1} = 1 \ GHz \gg f_3 \rightarrow$ It will not affect the phase margin

Feedforward Compensation

- In multistage amplifiers, usually there is one stage that acts as a bandwidth bottleneck by contributing a substantial amount of phase shift in the vicinity of the cross-over frequency f_x
- This technique creates a high-frequency bypass around the bottleneck stage to suppress its phase at f_x , thus improving ϕ_m



Feedforward Compensation

The bypass around the bottleneck stage is a high-pass function

$$h(jf) = \frac{jf/f_0}{1+jf/f_0}$$

The compensated open-loop gain is

$$a_{comp}(jf) = [a_1(jf) + h(jf)]a_2(jf)$$



• At low frequency: $|h(jf)| \ll |a_1(jf)|$

 $a_{comp}(jf) \cong a_1(jf)a_2(jf) = a(jf)$

The high low-frequency gain advantage of the uncompensated amplifier still hold.

• At high frequency:
$$|h(jf)| \gg |a_1(jf)|$$

 $a_{comp}(jf) \cong a_2(jf)$

The dynamics are controlled only by $a_2 \rightarrow$ Wider bandwidth & Lower phase shift

Summary of Internal Frequency Compensation

Dominant-pole compensation:



Additional Pole

- It creates an additional pole at sufficiently low frequency.
- It doesn't take advantage of the existing poles.
- It suffers from extremely low bandwidth.

Shunt-capacitance compensation:



- It rearranges the existing poles rather than creating an additional pole.
- It moves the first pole to sufficiently low frequency.
- The value of the shunt capacitance is extremely large \rightarrow Extra cost

Summary of Internal Frequency Compensation

Miller compensation:



- It takes advantage of Miller multiplicative effect of capacitors, thus requires moderate capacitance to move the first pole to sufficiently low frequency.
- It causes pole splitting, where the dominant pole is reduced and the first non-dominant pole is raised in frequency.

Summary of Internal Frequency Compensation

Pole-zero compensation:



- Similar to shunt-capacitance technique, a large capacitor is used to shift the first pole to sufficiently low frequency.
- A small resistance is used to create a zero that cancels the first nondominant pole
- Feedforward Compensation



• It places a high frequency bypass around the bottleneck stage that contributes the most phase shift in the vicinity of f_x

External Frequency Compensation

• How to stabilize the circuit by modifying its feedback factor β ?

- Reducing the Loop Gain
- Input-Lag Compensation
- Feedback-Lead Compensation

This method shifts $|1/\beta|$ curve upwards until it intercepts the |a| curve at $f = f_{\phi_m - 180^\circ}$, where ϕ_m is the desired phase margin.

The shift is obtained by connecting resistance R_c across the inputs.



 $\blacksquare R_c \text{ is chosen to achieve the desired phase margin } \phi_m:$



Reducing the Loop Gain

- Prices that we are paying for stability:
- Gain Error:

$$H_{CL} = \frac{1}{\beta} \frac{T}{1+T} = \frac{A_{ideal}}{1+1/T}$$

The presence of R_c reduces T, thus resulting in a larger gain error.

• DC Noise Gain:

$$H_{CL}(j0) \cong \frac{1}{\beta_0} = 1 + \frac{R_2}{R_1} + \frac{R_2}{R_c}$$

The presence of R_c causes an increased DC-noise gain which may result in an intolerable DC output error.

THERE'S NO FREE LUNCH !

Input-Lag Compensation

- The high DC-noise gain of the previous method can be overcome by placing a capacitance C_c in series with R_c .
 - High frequencies:
 - $\checkmark C_c$ is short.

 $\checkmark \frac{1}{|\beta|}$ curve is unchanged compared to the previous case.

• Low frequencies:



✓ $\frac{1}{|\beta|} = 1 + \frac{R_2}{R_1}$, we now have much higher DC loop gain & much lower DC output error.

R₂

[Franco]

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 \blacksquare R_c is chosen to achieve the desired phase margin ϕ_m :

$$\frac{1}{\beta_{\infty}} = 1 + \frac{R_2}{R_1} + \frac{R_2}{R_c} = \left| a(jf_{\phi_m - 180^\circ}) \right|$$
Then,
$$R_c = \frac{R_2}{\left| a(jf_{\phi_m - 180^\circ}) \right| - (1 + R_2/R_1)}$$

To avoid degrading ϕ_m , it is good practice to position the second breakpoint of $|1/\beta|$ curve a decade below $f_{\phi_m-180^\circ}$.

$$\frac{1}{2\pi C_c R_c} = \frac{1}{10} f_{\phi_m - 180^\circ}$$

Then,

$$C_c = \frac{5}{\pi R_c f_{\phi m^{-180^{\circ}}}}$$

Input-Lag Compensation

Advantage(s):

- \odot Lower DC-noise gain due to the presence of C_c .
- It allows for higher slew rate compared with internal compensation techniques: Op-amp is spared from having to charge/discharge internal compensation capacitance.

Disadvantage(s):

- \otimes Long settling tail because of the presence of pole-zero doublet of the feedback network ($|\beta|$).
- ☺ Increased high-frequency noise in the vicinity of the cross-over frequency.
- \otimes Low closed-loop differential input impedance (Z_d) which may cause high-frequency input loading

$$Z_d = z_d \| Z_c \quad , \quad Z_c = R_c + 1/sC_c \ll z_d$$

** z_d is the open loop input impedance of the Op-Amp.

Feedback-Lead Compensation

- This technique uses a feedback capacitance C_f to create phase lead in the feedback path.
- The phase lead is designed to be in the vicinity of the crossover frequency f_x which is were ϕ_m is boosted.



Feedback-Lead Compensation

Analysis:

$$\frac{1}{\beta} = 1 + \frac{\binom{R_2 \| Z_{C_f}}{R_1}}{R_1} = \left(1 + \frac{R_2}{R_1}\right) \frac{1 + jf/f_z}{1 + jf/f_p}$$

where,

$$f_p = \frac{1}{2\pi C_f R_2} \quad \text{,} \quad f_z = \left(1 + \frac{R_2}{R_1}\right) f_p$$

• The phase-lag provided by
$$\frac{1}{\beta(j\omega)}$$
 is maximum at $\sqrt{f_z f_p}$

• The optimum value of C_f , that maximizes the phase margin, is the one that makes this point at the crossover frequency.

$$f_x = \sqrt{f_z f_p} = f_p \sqrt{1 + \frac{R_2}{R_1}}$$

• The cross-over frequency can be obtained from

$$|a(jf_x)| = \frac{1}{|\beta(jf_x)|} = \sqrt{1 + \frac{R_2}{R_1}}$$

• Having f_x , the optimum C_f can be found

$$C_f = \frac{\sqrt{1 + R_2/R_1}}{2\pi R_2 f_x}$$


Feedback-Lead Compensation

How much phase margin can we get ?

• At the geometric mean of f_p and f_z , we have

$$\angle \left(\frac{1}{\beta}\right) = 90^{\circ} - 2\tan^{-1}\left(1 + \frac{R_2}{R_1}\right)$$

- The larger the value of $1 + R_2/R_1$, the greater the contribution of $1/\beta$ to the phase margin.
- E.g. $1 + R_2/R_1 = 10 \rightarrow \angle (1/\beta (jf_x)) = -55^\circ$ Thus,

$$\angle T(jf_x) = \angle a(jf_x) - \angle \left(\frac{1}{\beta(jf_x)}\right) = \angle a(jf_x) + 55^\circ$$

- The phase margin is improved by 55° due to feedback-lead compensation.



Feedback-Lead Compensation

Advantage(s):

- \bigcirc C_f helps to counteract the effect of the input stray capacitance C_n as we discussed beforehand.
- ③ It provides better filtering for internally generated noise.

Disadvantage(s):

③ It doesn't have the slew-rate advantage of the input-lag compensation.

These OpAmps are compensated for unconditional stability only when used with $1/\beta$ above a specified value

$$\frac{1}{\beta} \ge \left(\frac{1}{\beta}\right)_{min}$$

• They provide a constant GBP only for $|a| \ge (1/\beta)_{min}$

They offer higher GBP and slew rate.

Example

- The fully compensated LF356 OpAmp uses $C_c \cong 10 \, pF$ to provide $GBP = 5 \, MHz$ and $SR = 12 \, V/\mu s$ for any $|a| \ge 1 \, V/V$.
- The decompensated version of the same OpAmp, LF357, uses $C_c \cong 3 \, pF$ and provides $GBP = 20 \, MHz$ and $SR = 50 \, V/\mu s$ but only for any $|a| \ge 5 \, V/V$.